



## *Chapter 2: Motion and Recombination of Electrons and Holes*



## 2.1 Thermal Motion

Average electron kinetic energy =  $\frac{\text{total kinetic energy}}{\text{number of electrons}}$

$$= \frac{\int f(E)D(E)(E - E_c)dE}{\int f(E)D(E)dE}$$

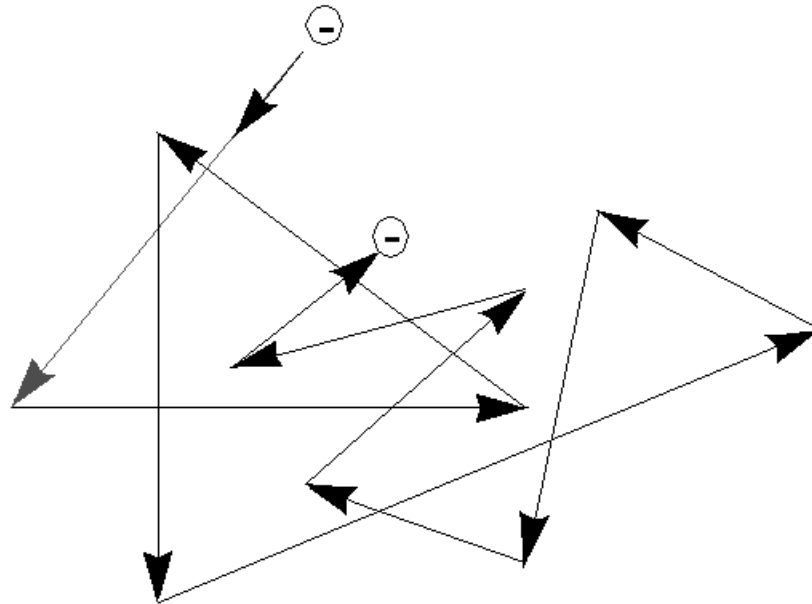
*Average electron or hole kinetic energy* =  $\frac{3}{2}kT = \frac{1}{2}mv_{th}^2$

$$v_{th} = \sqrt{\frac{3kT}{m_{eff}}} = \sqrt{\frac{3 \times 1.38 \times 10^{-23} \text{ JK}^{-1} \times 300\text{K}}{0.26 \times 9.1 \times 10^{-31} \text{ kg}}}$$

$$= 2.3 \times 10^5 \text{ m/s} = 2.3 \times 10^7 \text{ cm/s}$$



## 2.1 Thermal Motion

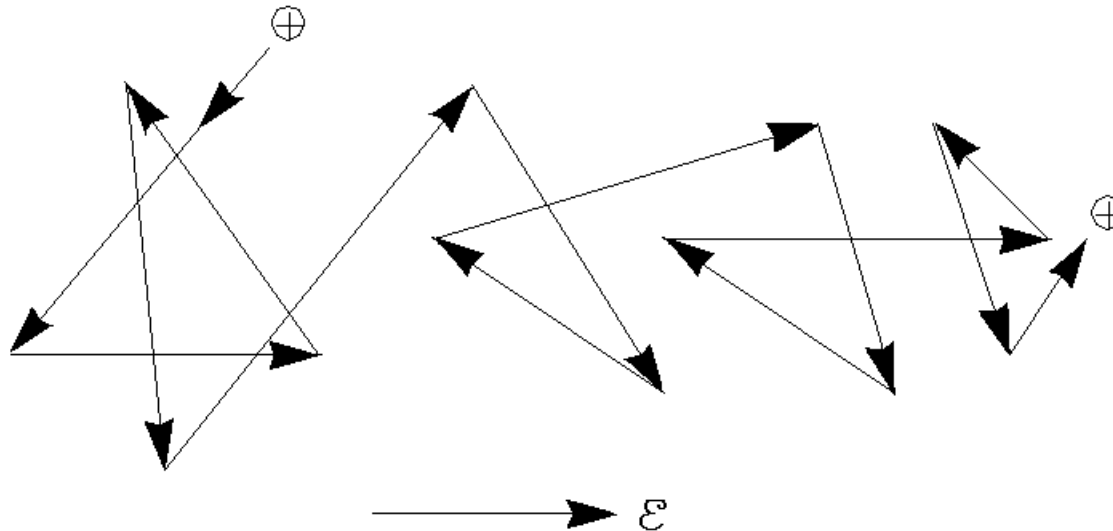


- Zig-zag motion is due to collisions or scattering with imperfections in the crystal.
- Net thermal velocity is zero.
- Mean time between collisions is  $\tau_m \sim 0.1\text{ps}$



## 2.2 Drift

### 2.2.1 Electron and Hole Mobilities



- *Drift* is the motion caused by an electric field.



## 2.2.1 Electron and Hole Mobilities

- The drift momentum gained between collisions is equal to the force,  $q\mathcal{E}$ , times the mean free time between collisions  $\tau_{mp}$ .

$$m_p v = q \mathcal{E} \tau_{mp}$$

$$v = \frac{q \mathcal{E} \tau_{mp}}{m_p}$$

$$v = \mu_p \mathcal{E}$$

$$\mu_p = \frac{q \tau_{mp}}{m_p}$$

$$v = -\mu_n \mathcal{E}$$

$$\mu_n = \frac{q \tau_{mn}}{m_n}$$

- $\mu_p$  is the hole mobility and  $\mu_n$  is the electron mobility



## 2.2.1 Electron and Hole Mobilities

$$v = \mu \varepsilon ; \quad \mu \text{ has the dimensions of } v/\varepsilon \quad \left[ \frac{\text{cm/s}}{\text{V/cm}} = \frac{\text{cm}^2}{\text{V} \cdot \text{s}} \right].$$

### *Electron and hole mobilities of selected semiconductors*

	<b>Si</b>	<b>Ge</b>	<b>GaAs</b>	<b>InAs</b>
$\mu_n$ (cm <sup>2</sup> /V·s)	1400	3900	8500	30000
$\mu_p$ (cm <sup>2</sup> /V·s)	470	1900	400	500

Based on the above table alone, which semiconductor and which carriers (electrons or holes) are attractive for applications in high-speed devices?



## Drift Velocity, Mean Free Time, Mean Free Path

**EXAMPLE:** Given  $\mu_p = 470 \text{ cm}^2/\text{V}\cdot\text{s}$ , what is the hole drift velocity at  $\mathcal{E} = 10^3 \text{ V/cm}$ ? What is  $\tau_{mp}$  and what is the distance traveled between collisions (called the **mean free path**)? Hint: When in doubt, use the MKS system of units.

**Solution:**  $v = \mu_p \mathcal{E} = 470 \text{ cm}^2/\text{V}\cdot\text{s} \times 10^3 \text{ V/cm} = 4.7 \times 10^5 \text{ cm/s}$

$$\begin{aligned}\tau_{mp} &= \mu_p m_p / q = 470 \text{ cm}^2/\text{V} \cdot \text{s} \times 0.39 \times 9.1 \times 10^{-31} \text{ kg} / 1.6 \times 10^{-19} \text{ C} \\ &= 0.047 \text{ m}^2/\text{V} \cdot \text{s} \times 2.2 \times 10^{-12} \text{ kg/C} = 1 \times 10^{-13} \text{ s} = 0.1 \text{ ps}\end{aligned}$$

$$\begin{aligned}\text{mean free path} &= \tau_{mh} v_{th} \sim 1 \times 10^{-13} \text{ s} \times 2.2 \times 10^7 \text{ cm/s} \\ &= 2.2 \times 10^{-6} \text{ cm} = 220 \text{ \AA} = 22 \text{ nm}\end{aligned}$$

This is smaller than the typical dimensions of devices, but getting close.



## 2.2.2 Mechanisms of Carrier Scattering

There are two main causes of carrier scattering:

1. Phonon Scattering
2. Ionized-Impurity (Coulombic) Scattering

**Phonon scattering** mobility decreases when temperature rises:

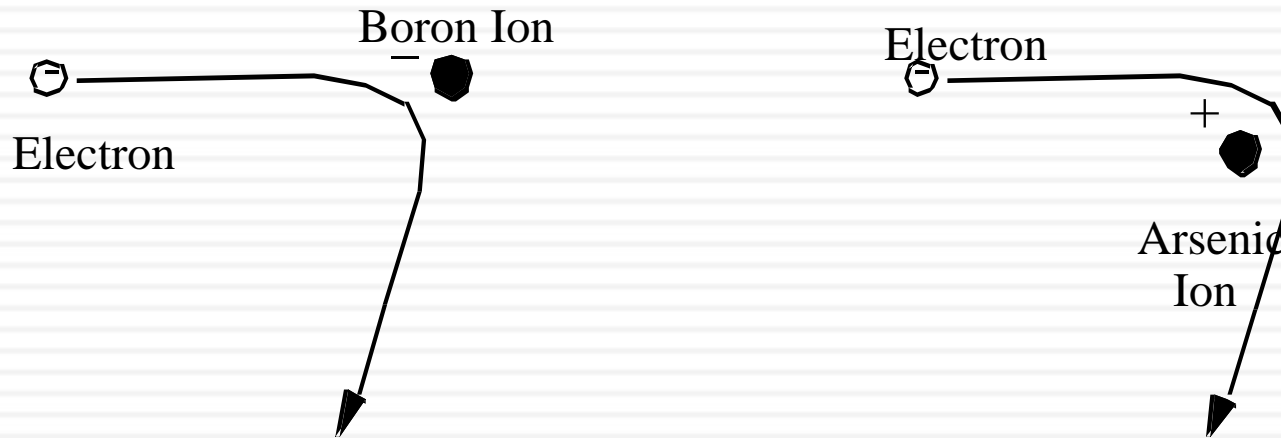
$$\mu_{phonon} \propto \tau_{phonon} \propto \frac{1}{\text{phonon density} \times \text{carrier thermal velocity}} \propto \frac{1}{T \times T^{1/2}} \propto T^{-3/2}$$

$\mu = q\tau/m$  (points to  $\mu_{phonon}$ )  
 $\propto T$  (points to phonon density)  
 $v_{th} \propto T^{1/2}$  (points to carrier thermal velocity)





# Impurity (Dopant)-Ion Scattering or Coulombic Scattering

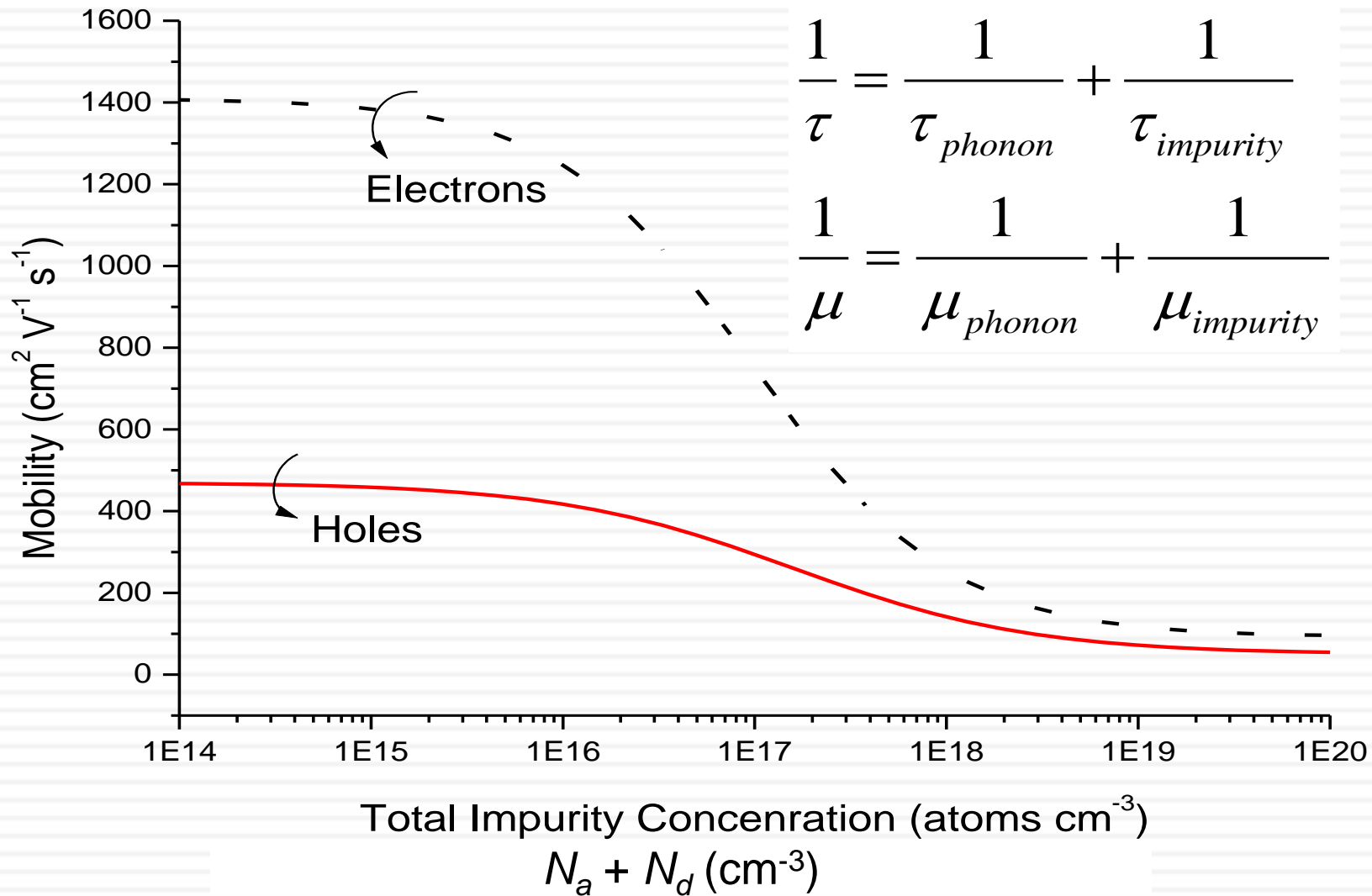


There is less change in the direction of travel if the electron zips by the ion at a higher speed.

$$\mu_{impurity} \propto \frac{v_{th}^3}{N_a + N_d} \propto \frac{T^{3/2}}{N_a + N_d}$$

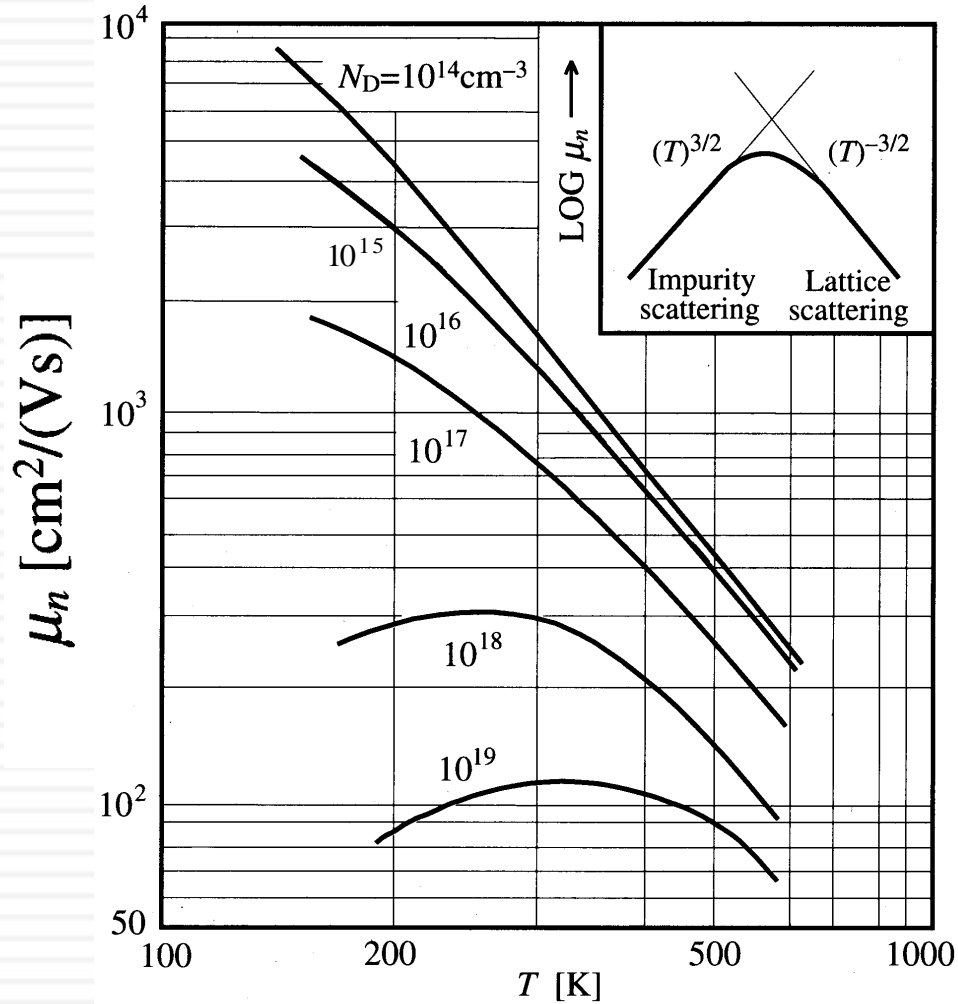


# Total Mobility





# Temperature Effect on Mobility



*Question:*

What  $N_d$  will make  $d\mu_n/dT = 0$  at room temperature?

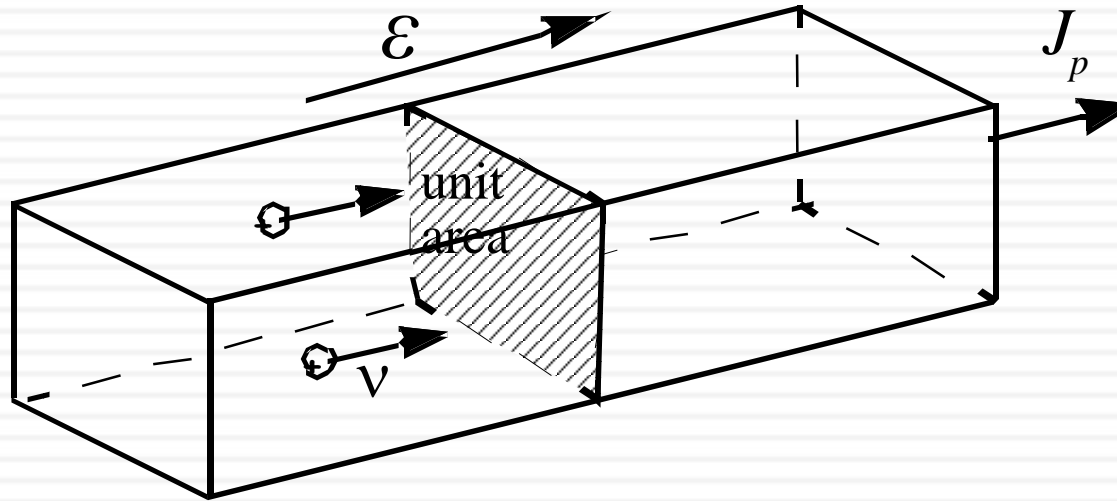


# Velocity Saturation

- When the kinetic energy of a carrier exceeds a critical value, it generates an optical photon and loses the kinetic energy.
- Therefore, the kinetic energy is capped at large  $\epsilon$ , and the velocity does not rise above a saturation velocity,  $v_{sat}$ .
- **Velocity saturation** has a deleterious effect on device speed as shown in Ch. 6.



## 2.2.3 Drift Current and Conductivity



Hole current density

$$J_p = qpv$$

A/cm<sup>2</sup> or C/cm<sup>2</sup>·sec

**EXAMPLE:** If  $p = 10^{15} \text{cm}^{-3}$  and  $v = 10^4 \text{cm/s}$ , then  
 $J_p = 1.6 \times 10^{-19} \text{C} \times 10^{15} \text{cm}^{-3} \times 10^4 \text{cm/s}$   
 $= 1.6 \text{C/s} \cdot \text{cm}^2 = 1.6 \text{A/cm}^2$



## 2.2.3 Drift Current and Conductivity

$$J_{p,drift} = qp v = qp \mu \varepsilon$$

$$J_{n,drift} = -qn v = qn \mu \varepsilon$$

$$J_{drift} = J_{n,drift} + J_{p,drift} = \sigma \varepsilon = (qn \mu_n + qp \mu_p) \varepsilon$$

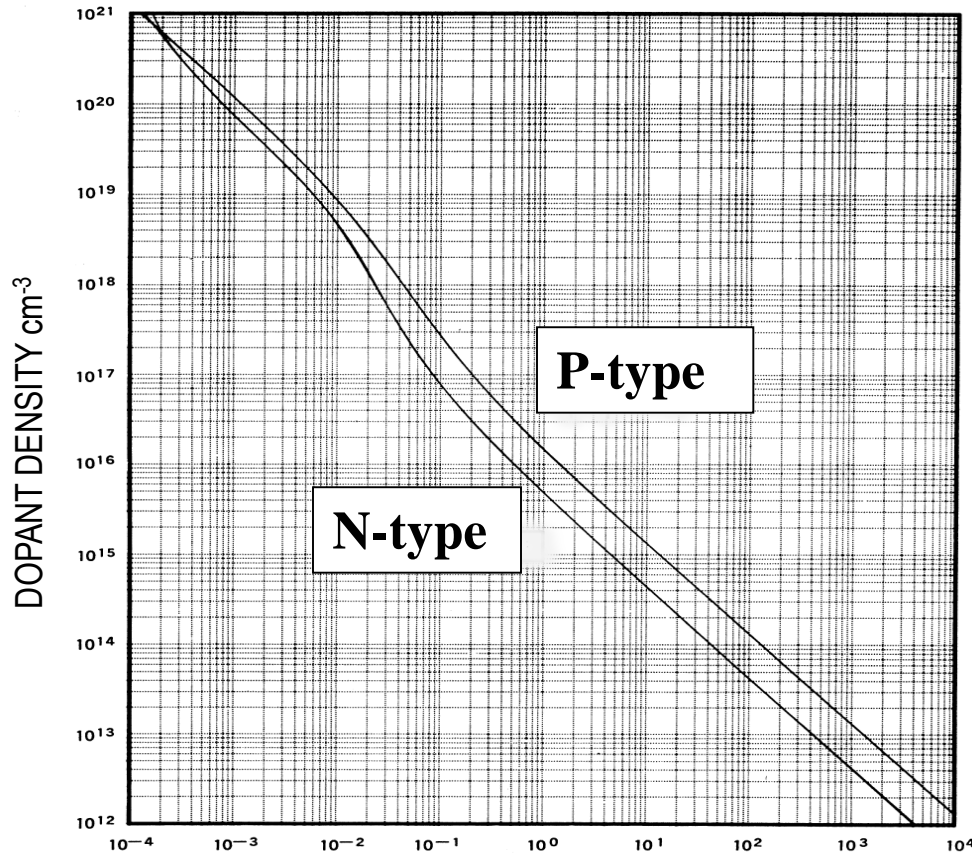
$\therefore$  **conductivity** (1/ohm-cm) of a semiconductor is

$$\sigma = qn \mu_n + qp \mu_p$$

$1/\sigma =$  is resistivity (ohm-cm)



# Relationship between Resistivity and Dopant Density



RESISTIVITY (Ω·cm)

$$\rho = 1/\sigma$$



## EXAMPLE: Temperature Dependence of Resistance

(a) What is the resistivity ( $\rho$ ) of silicon doped with  $10^{17} \text{cm}^{-3}$  of arsenic?

(b) What is the resistance ( $R$ ) of a piece of this silicon material  $1 \mu\text{m}$  long and  $0.1 \mu\text{m}^2$  in cross-sectional area?

### **Solution:**

(a) Using the  $N$ -type curve in the previous figure, we find that  $\rho = 0.084 \Omega\text{-cm}$ .

$$\begin{aligned} (b) R &= \rho L/A = 0.084 \Omega\text{-cm} \times 1 \mu\text{m} / 0.1 \mu\text{m}^2 \\ &= 0.084 \Omega\text{-cm} \times 10^{-4} \text{cm} / 10^{-10} \text{cm}^2 \\ &= 8.4 \times 10^{-4} \Omega \end{aligned}$$





## EXAMPLE: Temperature Dependence of Resistance

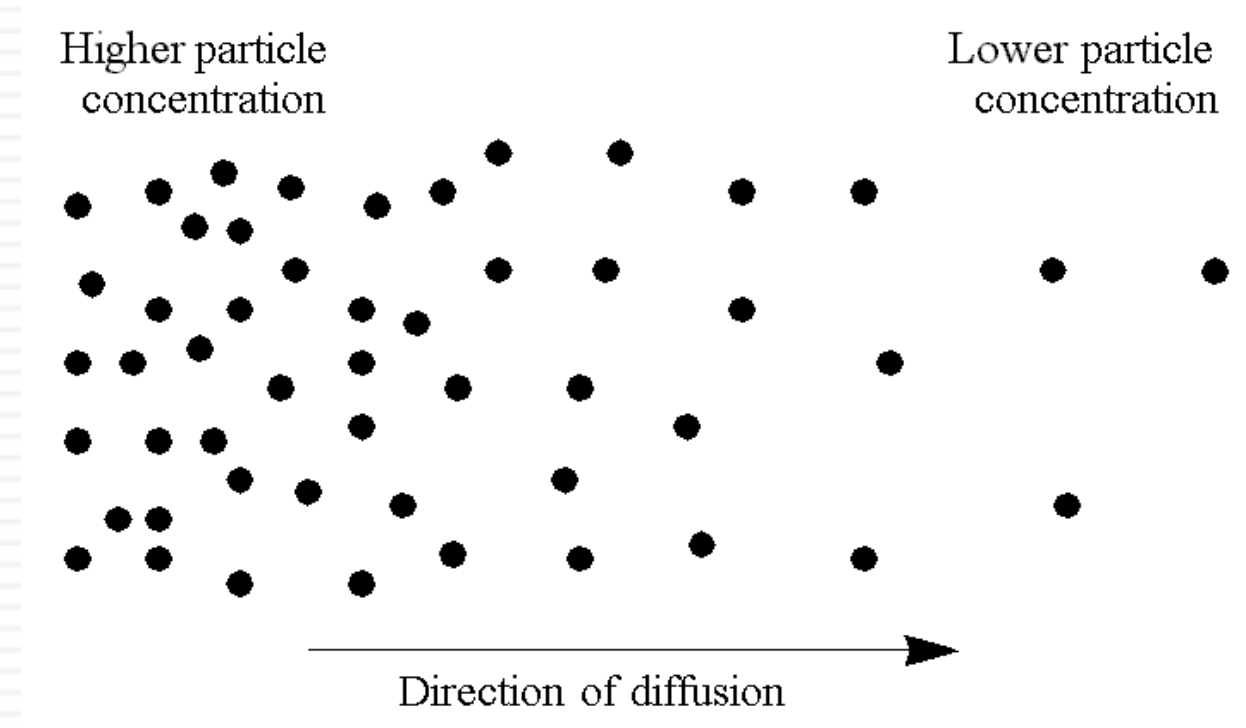
By what factor will  $R$  increase or decrease from  $T=300$  K to  $T=400$  K?

**Solution:** The temperature dependent factor in  $\sigma$  (and therefore  $\rho$ ) is  $\mu_n$ . From the mobility vs. temperature curve for  $10^{17}\text{cm}^{-3}$ , we find that  $\mu_n$  decreases from 770 at 300K to 400 at 400K. As a result,  $R$  **increases** by

$$\frac{770}{400} = 1.93$$



## 2.3 Diffusion Current



Particles diffuse from a higher-concentration location to a lower-concentration location.

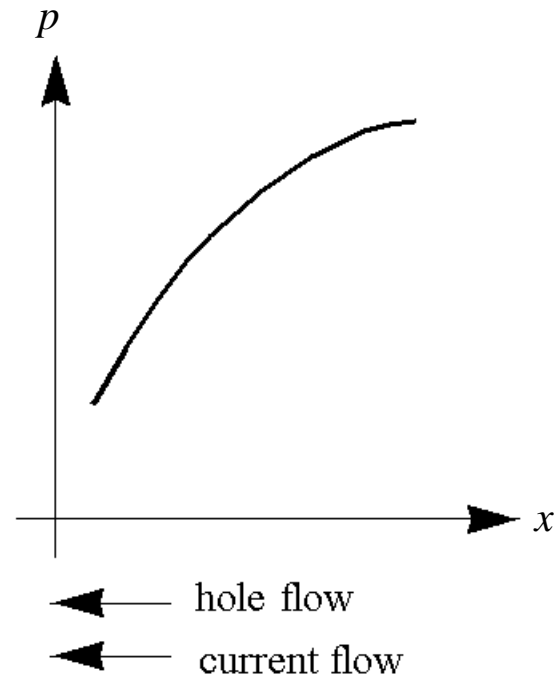
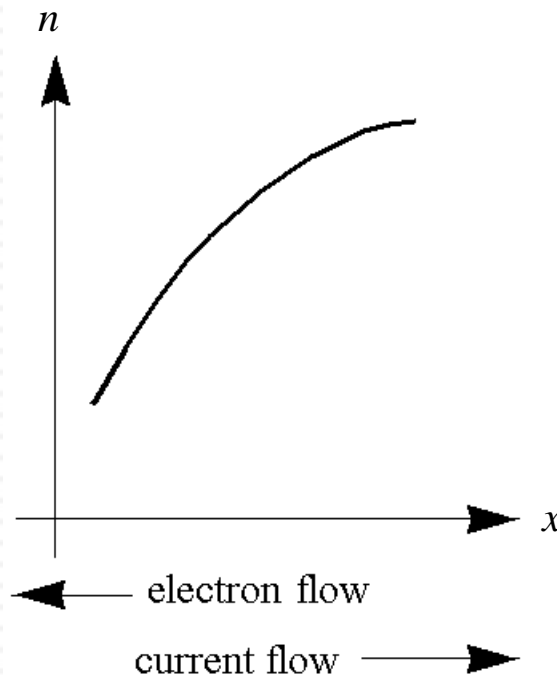


## 2.3 Diffusion Current

$$J_{n,diffusion} = qD_n \frac{dn}{dx}$$

$$J_{p,diffusion} = -qD_p \frac{dp}{dx}$$

$D$  is called the diffusion constant. Signs explained:





# Total Current – Review of Four Current Components

$$J_{TOTAL} = J_n + J_p$$

$$J_n = J_{n,drift} + J_{n,diffusion} = qn\mu_n \varepsilon + qD_n \frac{dn}{dx}$$

$$J_p = J_{p,drift} + J_{p,diffusion} = qp\mu_p \varepsilon - qD_p \frac{dp}{dx}$$



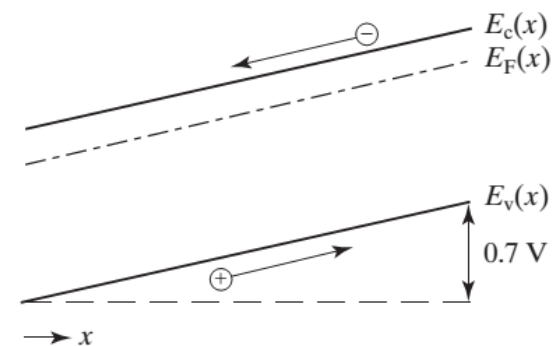
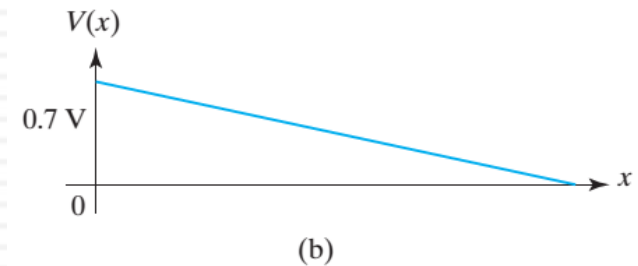
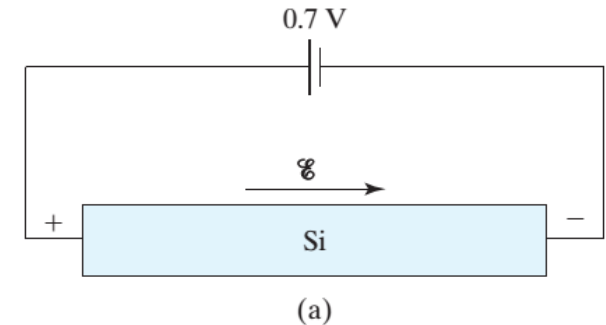
## 2.4 Relation Between the Energy Diagram and $V$ , $\mathcal{E}$

- When a voltage is applied across a piece of semiconductor as shown in Figure, it alters the band diagram
- A positive voltage raises the potential energy of a positive charge and lowers the energy of a negative charge

*$E_c$  and  $E_v$  vary in the opposite direction from the voltage. That is,  $E_c$  and  $E_v$  are higher where the voltage is lower.*

$$E_c(x) = \text{constant} - qV(x)$$

$$\mathcal{E}(x) = -\frac{dV}{dx} = \frac{1}{q} \frac{dE_c}{dx} = \frac{1}{q} \frac{dE_v}{dx}$$



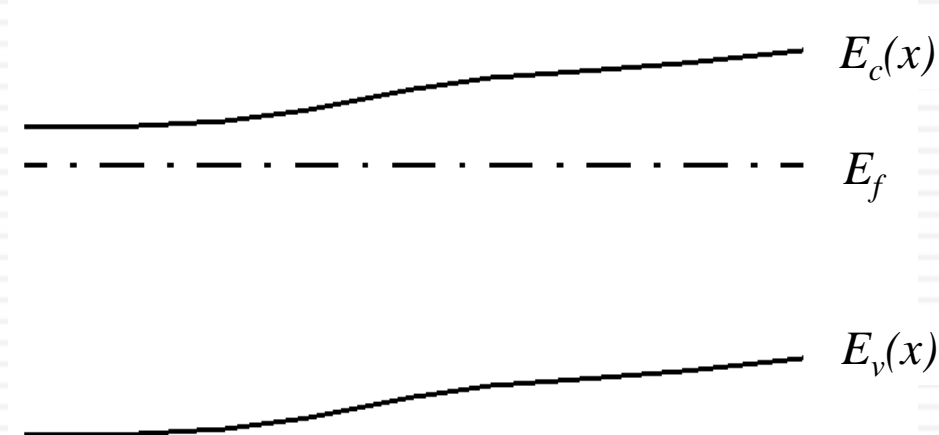


## 2.5 Einstein Relationship between $D$ and $\mu$

Consider a piece of non-uniformly doped semiconductor.

N-type semiconductor

Decreasing donor concentration



$$n = N_c e^{-(E_c - E_f)/kT}$$

$$\frac{dn}{dx} = -\frac{N_c}{kT} e^{-(E_c - E_f)/kT} \frac{dE_c}{dx}$$

$$= -\frac{n}{kT} \frac{dE_c}{dx}$$

$$= -\frac{n}{kT} q \mathcal{E}$$



## 2.5 Einstein Relationship between $D$ and $\mu$

$$\frac{dn}{dx} = -\frac{n}{kT} q\mathcal{E}$$

$$J_n = qn\mu_n\mathcal{E} + qD_n \frac{dn}{dx} = 0 \quad \text{at equilibrium.}$$

$$0 = qn\mu_n\mathcal{E} - qn \frac{qD_n}{kT} \mathcal{E}$$

$$D_n = \frac{kT}{q} \mu_n$$

Similarly,

$$D_p = \frac{kT}{q} \mu_p$$

*These are known as the **Einstein relationship**.*



## EXAMPLE: Diffusion Constant

What is the hole diffusion constant in a piece of silicon with  $\mu_p = 410 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  ?

**Solution:**

$$D_p = \left( \frac{kT}{q} \right) \mu_p = (26 \text{ mV}) \cdot 410 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1} = 11 \text{ cm}^2 / \text{s}$$

**Remember:**  $kT/q = 26 \text{ mV}$  at room temperature.





## 2.6 Electron-Hole Recombination

- The equilibrium carrier concentrations are denoted with  $n_0$  and  $p_0$ .
- The total electron and hole concentrations can be different from  $n_0$  and  $p_0$ .
- The differences are called the **excess carrier concentrations**  $n'$  and  $p'$ .

$$\begin{aligned}n &\equiv n_0 + n' \\ p &\equiv p_0 + p'\end{aligned}$$



# Charge Neutrality

- Charge neutrality is satisfied at equilibrium ( $n' = p' = 0$ ).
- When a non-zero  $n'$  is present, an equal  $p'$  may be assumed to be present to maintain charge equality and vice-versa.
- If charge neutrality is not satisfied, the net charge will attract or repel the (majority) carriers through the drift current until neutrality is restored.

$$n' = p'$$



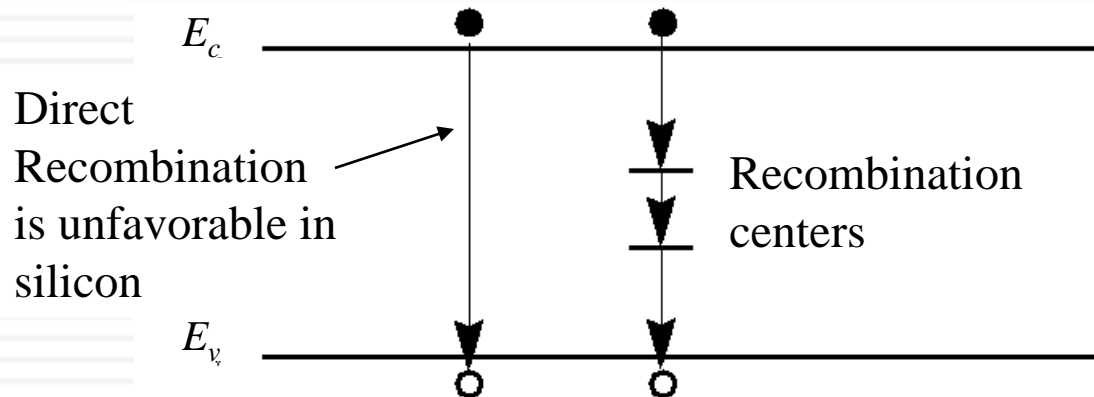
## Recombination Lifetime

- Assume light generates  $n'$  and  $p'$ . If the light is suddenly turned off,  $n'$  and  $p'$  decay with time until they become zero.
- The process of decay is called **recombination**.
- The time constant of decay is the **recombination time** or **carrier lifetime**,  $\tau$ .
- Recombination is nature's way of restoring equilibrium ( $n' = p' = 0$ ).



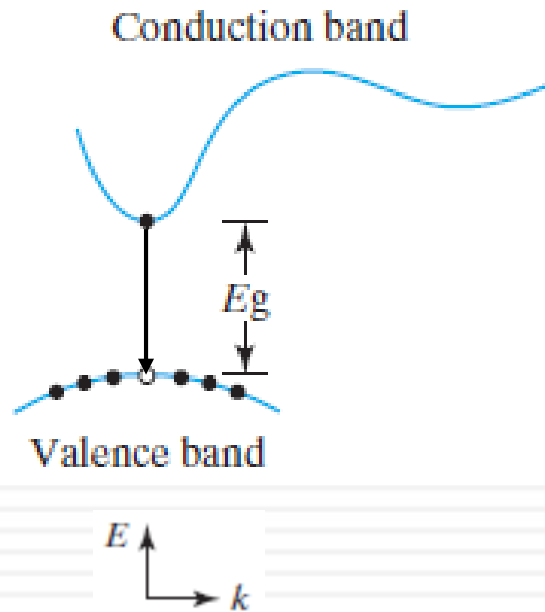
# Recombination Lifetime

- $\tau$  ranges from 1ns to 1ms in Si and depends on the density of metal impurities (contaminants) such as Au and Pt.
- These **deep traps** capture electrons and holes to facilitate recombination and are called **recombination centers**.



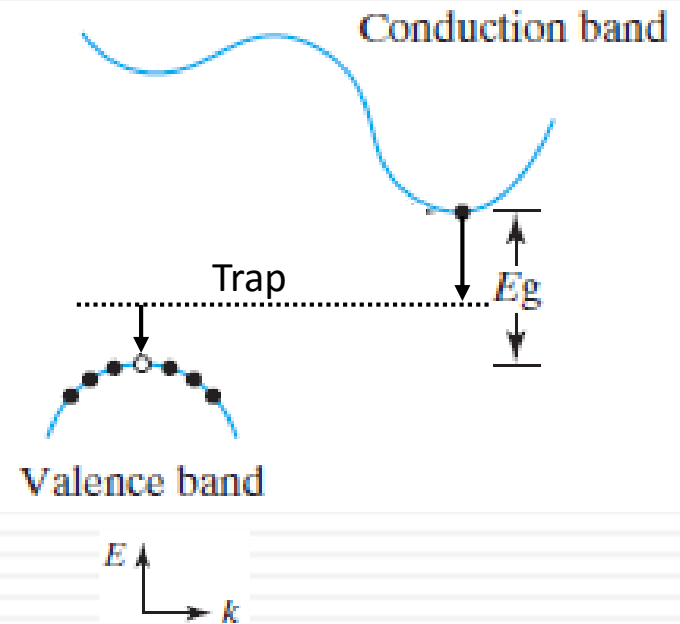


# Direct and Indirect Band Gap



Direct band gap  
Example: GaAs

Direct recombination is efficient  
as  $k$  conservation is satisfied.



Indirect band gap  
Example: Si

Direct recombination is rare as  $k$   
conservation is not satisfied



## Rate of recombination ( $s^{-1}cm^{-3}$ )

$$\frac{dn'}{dt} = -\frac{n'}{\tau}$$

$$n' = p'$$

$$\frac{dn'}{dt} = -\frac{n'}{\tau} = -\frac{p'}{\tau} = \frac{dp'}{dt}$$



## EXAMPLE: Photoconductors

- A bar of Si is doped with boron at  $10^{15}\text{cm}^{-3}$ . It is exposed to light such that electron-hole pairs are generated throughout the volume of the bar at the rate of  $10^{20}/\text{s}\cdot\text{cm}^3$ . The recombination lifetime is  $10\mu\text{s}$ . What are (a)  $p_0$ , (b)  $n_0$ , (c)  $p'$ , (d)  $n'$ , (e)  $p$ , (f)  $n$ , and (g) the  $np$  product?



## EXAMPLE: Photoconductors

### Solution:

(a) What is  $p_o$ ?

$$p_o = N_a = 10^{15} \text{ cm}^{-3}$$

(b) What is  $n_o$ ?

$$n_o = n_i^2/p_o = 10^5 \text{ cm}^{-3}$$

(c) What is  $p'$ ?

*In steady-state, the rate of generation is equal to the rate of recombination.*

$$10^{20}/\text{s-cm}^3 = p'/\tau$$

$$\therefore p' = 10^{20}/\text{s-cm}^3 \cdot 10^{-5}\text{s} = 10^{15} \text{ cm}^{-3}$$





## EXAMPLE: Photoconductors

(d) What is  $n'$ ?

$$n' = p' = 10^{15} \text{ cm}^{-3}$$

(e) What is  $p$ ?

$$p = p_0 + p' = 10^{15} \text{ cm}^{-3} + 10^{15} \text{ cm}^{-3} = 2 \times 10^{15} \text{ cm}^{-3}$$

(f) What is  $n$ ?

$$n = n_0 + n' = 10^5 \text{ cm}^{-3} + 10^{15} \text{ cm}^{-3} \sim 10^{15} \text{ cm}^{-3} \text{ since } n_0 \ll n'$$

(g) What is  $np$ ?

$$np \sim 2 \times 10^{15} \text{ cm}^{-3} \cdot 10^{15} \text{ cm}^{-3} = 2 \times 10^{30} \text{ cm}^{-6} \gg n_i^2 = 10^{20} \text{ cm}^{-6}.$$

The  $np$  product can be very different from  $n_i^2$ .



## 2.7 Thermal Generation

If  $n'$  is negative, there are fewer electrons than the equilibrium value.

As a result, there is a net rate of ***thermal generation*** at the rate of  $|n'|/\tau$ .



## 2.8 Quasi-equilibrium and Quasi-Fermi Levels

- Whenever  $n' = p' \neq 0$ ,  $np \neq n_i^2$ . We would like to preserve and use the simple relations:

$$n = N_c e^{-(E_c - E_f)/kT}$$

$$p = N_v e^{-(E_f - E_v)/kT}$$

- But these equations lead to  $np = n_i^2$ . The solution is to introduce two **quasi-Fermi levels**  $E_{fn}$  and  $E_{fp}$  such that

$$n = N_c e^{-(E_c - E_{fn})/kT}$$

$$p = N_v e^{-(E_{fp} - E_v)/kT}$$

Even when electrons and holes are not at equilibrium, *within each group* the carriers can be at equilibrium. Electrons are closely linked to other electrons but only loosely to holes.



## EXAMPLE: Quasi-Fermi Levels and Low-Level Injection

Consider a Si sample with  $N_d = 10^{17} \text{ cm}^{-3}$  and  $n' = p' = 10^{15} \text{ cm}^{-3}$ .

(a) Find  $E_f$ .

$$n = N_d = 10^{17} \text{ cm}^{-3} = N_c \exp[-(E_c - E_f)/kT]$$

$$\therefore E_c - E_f = 0.15 \text{ eV. } (E_f \text{ is below } E_c \text{ by } 0.15 \text{ eV.})$$

Note:  $n'$  and  $p'$  are much less than the majority carrier concentration. This condition is called **low-level injection**.